

## Raman spectra investigation of InAlGaN quaternary alloys grown by metalorganic chemical vapor deposition

S. Y. Hu, Y. C. Lee, Z. C. Feng, and Y. H. Weng

Citation: *J. Appl. Phys.* **112**, 063111 (2012); doi: 10.1063/1.4752420

View online: <http://dx.doi.org/10.1063/1.4752420>

View Table of Contents: <http://jap.aip.org/resource/1/JAPIAU/v112/i6>

Published by the [American Institute of Physics](#).

---

### Related Articles

Temperature dependence of Mg-H local vibrational modes in heavily doped InN:Mg  
*J. Appl. Phys.* **112**, 053528 (2012)

Lamb wave band gaps in a homogenous plate with periodic tapered surface  
*J. Appl. Phys.* **112**, 054503 (2012)

Selectively probing vibrations in a plasmonic supracrystal  
*Appl. Phys. Lett.* **101**, 101903 (2012)

Effect of ternary mixed crystals on interface optical phonons in wurtzite  $\text{In}_x\text{Ga}_{1-x}\text{N}/\text{GaN}$  quantum wells  
*J. Appl. Phys.* **112**, 053704 (2012)

Role of phonon dispersion in studying phonon mean free paths in skutterudites  
*J. Appl. Phys.* **112**, 044305 (2012)

---

### Additional information on *J. Appl. Phys.*

Journal Homepage: <http://jap.aip.org/>

Journal Information: [http://jap.aip.org/about/about\\_the\\_journal](http://jap.aip.org/about/about_the_journal)

Top downloads: [http://jap.aip.org/features/most\\_downloaded](http://jap.aip.org/features/most_downloaded)

Information for Authors: <http://jap.aip.org/authors>

## ADVERTISEMENT



**AIP Advances**

Special Topic Section:  
**PHYSICS OF CANCER**

Why cancer? Why physics? [View Articles Now](#)

## Raman spectra investigation of InAlGaN quaternary alloys grown by metalorganic chemical vapor deposition

S. Y. Hu,<sup>1,a)</sup> Y. C. Lee,<sup>2</sup> Z. C. Feng,<sup>3</sup> and Y. H. Weng<sup>4</sup>

<sup>1</sup>Department of Electrical Engineering, Tungfang Design University, Hunei, Kaohsiung 82941, Taiwan

<sup>2</sup>Department of Electronic Engineering, Tungnan University, Shen-Keng, Taipei 22202, Taiwan

<sup>3</sup>Graduate Institute of Electro-Optical Engineering, National Taiwan University, Taipei 10617, Taiwan

<sup>4</sup>Department of Electrical Engineering, National Taiwan Ocean University, Keelung 20224, Taiwan

(Received 29 May 2012; accepted 15 August 2012; published online 19 September 2012)

Raman analysis of the  $A_1$  (LO) mode of AlGa<sub>N</sub>-like phonons of the InAlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures in the composition range  $1.38\% \leq \text{In} \leq 2.73\%$  and  $8.01\% \leq \text{Al} \leq 13.97\%$  is presented. The line shape of  $A_1$  (LO) mode of AlGa<sub>N</sub>-like phonons was observed to exhibit a significant asymmetry and Raman linewidth toward the lower energy side. The spatial correlation model is discussed and is shown to account the line shape. The spatial correlation model calculations also indicate the lack of a long-range order in the higher Al to In ratio of InAlGa<sub>N</sub>/Ga<sub>N</sub> alloys. These results were confirmed by x-ray diffraction and the correlation length  $L$  decreases as the increasing of Al to In ratio corresponding to the absence of the long-range order in the alloy. The Raman linewidth of the AlGa<sub>N</sub>-like  $A_1$  (LO) mode was found to exhibit a maximum at the higher Al to In ratio indicative of a random disordered alloy system. © 2012 American Institute of Physics. [<http://dx.doi.org/10.1063/1.4752420>]

### I. INTRODUCTION

The importance of InAlGa<sub>N</sub> system of quaternary alloys offers the opportunities to modify their fundamental physical properties over a wide range by purposely varying In or Al compositions in InAlGa<sub>N</sub> to change the band gap while keeping lattice matched with Ga<sub>N</sub>, which can be used to reduce dislocation density and piezoelectric field.<sup>1–3</sup> Therefore, InAlGa<sub>N</sub> quaternary alloys have been successfully introduced to be good for luminescence enhancement, especially in the ultraviolet (UV) regions.<sup>1–3</sup> In spite of the significant progress in the preparation of these quaternary alloys of InAlGa<sub>N</sub> achieved in recent years by adjusting the alloy compositions, many of their basic characteristics still remain uncertain.<sup>4–6</sup> Furthermore, there are few papers focused on the typical feature of alloys for a disorder in atomic distributions over lattice sites to bring out the properties of electron states or lattice dynamics.<sup>4–6</sup> Particularly, it is true for the phonon spectrum because such an analysis can provide not only insight into the atomic structures but also bonding properties to direct a better understanding of the electronic properties and applications in the InAlGa<sub>N</sub> quaternary alloys.<sup>4–6</sup> Raman scattering has been used to study the physical properties of InAlGa<sub>N</sub> quaternary alloys and the Raman spectra show difference of various phonon modes with compositional disorder, including the phonon frequency shift, changes of linewidth and asymmetry, and coming out of disorder-activated modes.<sup>4–6</sup> Recently, considerable attention has been devoted to the discussions of the phonons related to the InGa<sub>N</sub>-like  $A_1$  (LO) modes or AlGa<sub>N</sub>-like  $A_1$  (LO) modes in the InAlGa<sub>N</sub> quaternary alloys.<sup>4–6</sup> However, the relation between alloy disorder and the line shape (linewidth and asymmetry) of  $A_1$  (LO) mode of AlGa<sub>N</sub>-like pho-

nons in the InAlGa<sub>N</sub> quaternary alloys has not been fully studied.

In this paper, a detailed investigation on the influence of the Raman spectra caused by alloy compositional fluctuations (ACFs) in the different quaternary alloys of InAlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures will be provided. Moreover, we pay attention to the Ga<sub>N</sub>  $A_1$  (LO) mode and AlGa<sub>N</sub>-like  $A_1$  (LO) mode in the Raman shift range of 675–800  $\text{cm}^{-1}$  at room temperature. The linewidth and asymmetry of Raman line shape in the AlGa<sub>N</sub>-like  $A_1$  (LO) mode can be investigated in terms of the spatial correlation (SC) model regarding to the finite correlation length of a propagating phonon due to the ACFs.<sup>7</sup>

### II. EXPERIMENTAL

Quaternary heterostructure samples of the alloy system In-Al-Ga-N with indium (In) content from 1.38% to 2.73% and aluminum (Al) content from 8.01% to 13.97% were grown on the c-plane sapphire substrates by metalorganic chemical vapor deposition (MOCVD) method. Each of the produced quaternary alloys were grown on top of the AlN buffer layer (25 nm) and consisted of one InAlGa<sub>N</sub> layer (0.23  $\mu\text{m}$ ) with a Ga<sub>N</sub> layer (0.25  $\mu\text{m}$ ) in between. The growth temperature of 1050 °C and pressure of 300 Torr were applied for the growth of underlying Ga<sub>N</sub> epilayer. InAlGa<sub>N</sub> layers were grown at 780 °C, while In and Al compositions were controlled by varying the flow rates of trimethylindium (TMIn) and trimethylaluminum (TMAl), respectively. The alloy compositions were determined by x-ray photoelectron spectroscopy (XPS) and are listed in Table I for all of the three samples. The structural properties of the InAlGa<sub>N</sub>/Ga<sub>N</sub> heterostructure quaternary alloys were analyzed by high resolution x-ray diffraction (HRXRD). The Raman spectra were measured at room temperature by using the Renishaw System (inVia Raman microscope) while the spectra were excited using the 514.5 nm line of an Ar<sup>+</sup> laser at an incident power

<sup>a)</sup>Author to whom correspondence should be addressed. Electronic mail: shenghu2729@yahoo.com.

TABLE I. Sample number, composition, ratio of Al to In, the values of linewidth  $\Gamma$ , and correlation length  $L$  for the AlGaIn-like  $A_1$  (LO) phonon modes in our InAlGaIn/GaN heterostructure quaternary alloys.

Sample	In (%)	Al (%)	Al to In	$\Gamma$ ( $\text{cm}^{-1}$ )	$L$ ( $\text{\AA}$ )
M12	1.38	10.52	7.6	15.0	62
M14	2.73	13.97	5.1	14.0	66
M28	1.64	8.01	4.9	13.0	69

of 10 mW with the experimental error  $\pm 1 \text{ cm}^{-1}$  for linewidth and  $\pm 0.5 \text{ cm}^{-1}$  for peak position, respectively.

### III. RESULTS AND DISCUSSION

Room temperature Raman spectra of the different InAlGaIn samples over the range of  $450\text{--}800 \text{ cm}^{-1}$  are illustrated in Fig. 1. For the InAlGaIn layer grown on GaN, those Raman peaks have shown that there are three resolved phonon structures observed in each Raman spectrum of InAlGaIn quaternary alloys and the assignments of our Raman modes are based on the previous Raman studies of InAlGaIn quaternary alloys.<sup>2-6</sup> In (a)–(c) of Fig. 1, one peak position located in the region of  $550\text{--}575 \text{ cm}^{-1}$  corresponds to the GaN  $E_2(\text{high})$  mode and GaN  $A_1$  (LO) mode

$A_1$  (LO) mode is located at around  $735 \text{ cm}^{-1}$ , respectively. Furthermore, from  $700$  to  $800 \text{ cm}^{-1}$ , the AlGaIn-like  $A_1$  (LO) mode is located at around  $750 \text{ cm}^{-1}$  described by the Gaussian fit (dotted line) of the Raman line shape and the spectra are displayed in (d)–(f) of Fig. 1. The  $A_1$  (LO) mode of InGaN-like, expected at around  $700 \text{ cm}^{-1}$  is not detectable, probably it is because the In composition is much less than Al composition in our quaternary samples.<sup>2-6</sup> Then, it is seen that most of the Raman peaks from our samples correspond well to those of the other InAlGaIn quaternary alloys,<sup>2-6</sup> the Raman spectra of our quaternary samples have shown their particular characteristics to verify that both of the  $E_2(\text{high})$  mode of GaN and  $A_1$  (LO) mode of GaN are from the substrates, while the  $A_1$  (LO) modes of AlGaIn-like are similar to the Al-rich  $A_1$  (LO) phonon lines from the earlier reports.<sup>2-6</sup>

Figure 2 shows the room temperature Raman spectra of the AlGaIn-like  $A_1$  (LO) modes from our quaternary samples M12, M14, and M28 in the range of  $675\text{--}800 \text{ cm}^{-1}$ , respectively. As displayed in Fig. 2, the spectra line shape for all three samples reveal asymmetric broadening (Raman linewidth) and a shift to lower frequency of the AlGaIn-like  $A_1$  (LO) mode with increasing the ratio of Al to In. To the best of our knowledge, both topological and structural

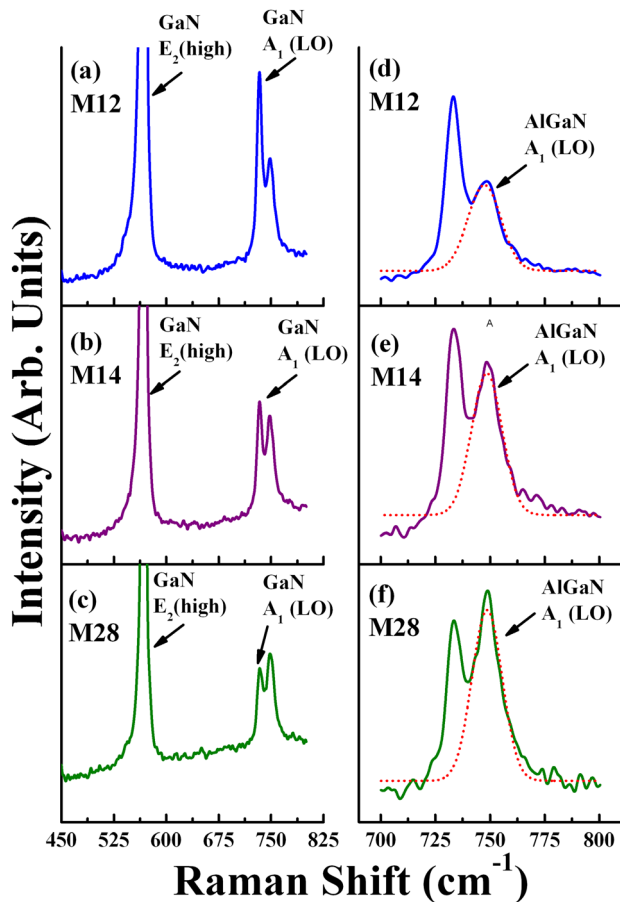


FIG. 1. Room temperature Raman spectra of the different InAlGaIn samples over the range of  $450\text{--}800 \text{ cm}^{-1}$  are illustrated. In (a)–(c), one peak position located in the region of  $550\text{--}575 \text{ cm}^{-1}$  corresponds to the GaN  $E_2(\text{high})$  mode and GaN  $A_1$  (LO) mode located at around  $735 \text{ cm}^{-1}$ , respectively. In (d)–(f), from  $700$  to  $800 \text{ cm}^{-1}$ , the AlGaIn-like  $A_1$  (LO) mode is located at around  $750 \text{ cm}^{-1}$  described by the Gaussian fit (dotted line) of the Raman line shape.

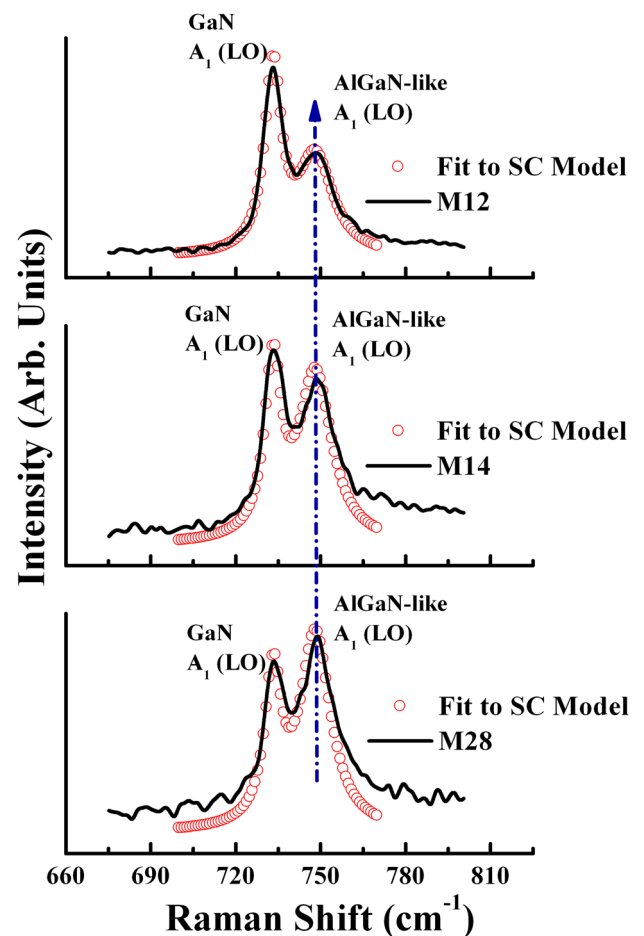


FIG. 2. Room temperature Raman spectra of the AlGaIn-like  $A_1$  (LO) modes from our quaternary samples M12, M14, and M28 are in the range of  $675\text{--}800 \text{ cm}^{-1}$ , respectively. The spectra line shape for all three samples reveal asymmetric Raman linewidth and a shift to lower frequency of the AlGaIn-like  $A_1$  (LO) mode with increasing the ratio of Al to In.

disorder may result from the process of atom substitution in alloy semiconductors. For allowed phonons, these disorders result basically in the breaking of the translational symmetry, leading to the contribution of  $q \neq 0$  phonons to the Raman line shape, corresponding to the so-called finite-size effects.<sup>7–12</sup> This extensively drives the line shape asymmetry, which can be analyzed in terms of the SC model regarding to the finite correlation length of a propagating phonon due to the ACFs.<sup>7–12</sup>

For an ideal crystal, because of the momentum conservation, only phonons at the center of the Brillouin zone ( $q=0$ ) can be observed by Raman scattering. As the crystal is alloying, the phonons can be confined in space owing to the potential fluctuations of the alloy disorder, which give rise to a relaxation of the  $q=0$  selection rule in Raman scattering.<sup>7–12</sup> Then, the spatial correlation length of phonon in alloys becomes finite. The finite phonon mode will lead to the broadening and asymmetry of the Raman line shape. With the SC model, we can evaluate the asymmetric broadening of Raman scattering by the theoretical calculation. The Raman peak of AlGaIn-like  $A_1$  (LO) mode in Fig. 2 shows an asymmetric line shape, which can be described by the SC model.<sup>7–12</sup>

Based on the SC model, the Gaussian spatial correlation function  $\exp(-2r^2/L^2)$ , where  $L$  is the phonon correlation length, has been effectively used to account for  $q$ -vector relaxation related to finite-size effects and structural disorder.<sup>7–12</sup> Then, the Raman intensity  $I(\omega)$  at a frequency  $\omega$  can be written as follows:

$$I(\omega) \propto \int \frac{\exp(-q^2 L^2/4)}{[\omega - \omega(q)]^2 + (\Gamma_0/2)^2} d^3 q, \quad (1)$$

where  $q$  is the reduced wave vector in the unit of  $2\pi/a$ ,  $a$  is the lattice constant, and  $\Gamma_0$  is the linewidth of the material which is GaN and AlGaIn in this study.<sup>7–12</sup> In our case, the Raman linewidth is around  $7 \text{ cm}^{-1}$  for GaN  $A_1$  (LO) mode and  $13\text{--}15 \text{ cm}^{-1}$  for AlGaIn-like  $A_1$  (LO) mode, respectively. Here, we consider that the Gaussian distribution is associated with the finite size of thin film region and rewrite Eq. (1) as follows:<sup>10</sup>

$$I(\omega) \propto \int \frac{\exp(-q^2 L_1^2/16\pi^2)}{[\omega_1 - \omega_1(q)]^2 + (\Gamma_1/2)^2} d^3 q + \int \frac{\exp(-q^2 L_2^2/16\pi^2)}{[\omega_2 - \omega_2(q)]^2 + (\Gamma_2/2)^2} d^3 q. \quad (2)$$

For the dispersion  $\omega_1(q)$  and  $\omega_2(q)$ , we take the following analytic model relationship:<sup>7–12</sup>

$$\omega_1(q) = C_1 + D_1 \cos(\pi q) \quad \text{and} \quad \omega_2(q) = C_2 + D_2 \cos(\pi q). \quad (3)$$

According to the *ab initio* phonon-dispersion relations,  $C_1$  and  $D_1$  were calculated for the GaN  $A_1$  (LO) mode, while  $C_2$  and  $D_2$  were calculated for the AlGaIn-like  $A_1$  (LO) mode, respectively. Equations (2) and (3) show that if  $L$  is finite, the  $q$  selection rule may relax and additional transitions by

phonon with lower energy at  $q \neq 0$  may occur, which lead to the broadening and asymmetry of the Raman line shape. The opened circles in Fig. 2 exhibit the calculated results with the SC model, which are somewhat in good agreement with the experimental data. At present, we merely focus on the transformation of the correlation length  $L$  and the Raman linewidth  $\Gamma$  in AlGaIn-like  $A_1$  (LO) mode assigned from our samples, then the values of  $L$  and  $\Gamma$  can be evaluated from the Raman line shape fitting and are also listed in Table I.

From the results of Table I, as the ratio of Al to In increased, it is found that the correlation length  $L$  decreases but the Raman linewidth  $\Gamma$  increases. The correlation length  $L$  is determined from our calculations to be in the range of  $62\text{--}69 \text{ \AA}$ , a relatively small value suggesting a clear absence of long-range order in the alloy system.<sup>5</sup> In order to further investigate this assumption, each of our InAlGaIn quaternary alloys were characterized by the HRXRD measurements to carry out the (004)  $\omega - 2\theta$  scan. The HRXRD spectra are shown in Fig. 3 and are relatively comparable to our previous results.<sup>13</sup> In Fig. 3, the separated diffraction peaks are associated with GaN and InAlGaIn, respectively, which confirm that the quaternary InAlGaIn alloys were epitaxially grown on GaN. The alloy system of In-Al-Ga-N samples with the Al to In ratio of around 4.8, which are lattice matched to the underlying GaN layer, exhibited the best structural and optical properties, so only one GaN peak is observed for M28 in Fig. 3.<sup>14</sup> However, as the Al to In ratio increased more, the InAlGaIn diffraction peaks appear on the right side of GaN peak. When the Al to In ratio is increased further to be around 7.6, a more isolated InAlGaIn peak appears. The separated distance between GaN and InAlGaIn peaks in M12 (dotted line) is carried out to be the largest one, demonstrating the biggest lattice mismatch and strain existing in M12 among the three samples.<sup>14</sup> Therefore, the HRXRD confirms the Raman analysis that M12 in this study does not reveal considerable long-range ordering.

More studies related to the InAlGaIn/GaN heterostructures may be obtained from the Raman linewidth and

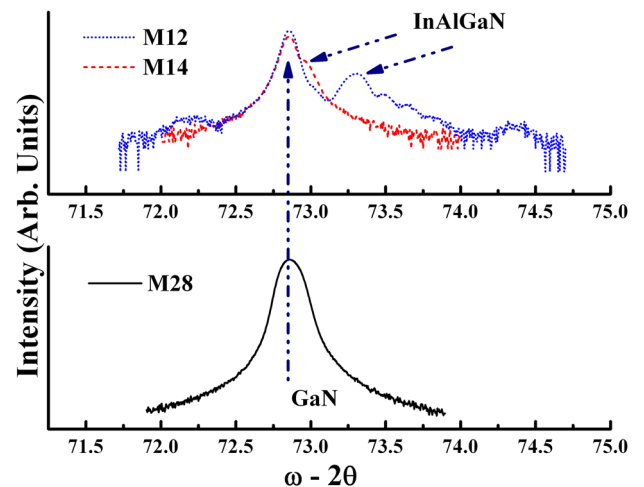


FIG. 3. Three samples of our InAlGaIn/GaN heterostructures were characterized by HRXRD to carry out the (004)  $\omega - 2\theta$  scan. The separated diffraction peaks are associated with GaN and InAlGaIn, respectively, which confirm that the quaternary InAlGaIn alloys were epitaxially grown on GaN.

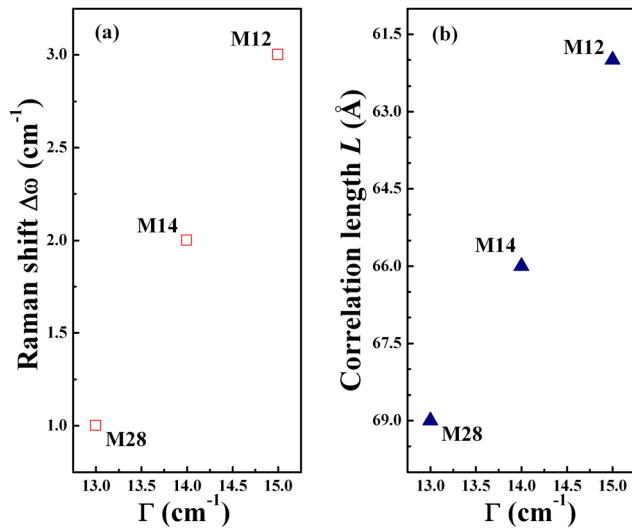


FIG. 4. The Raman shift  $\Delta\omega$  (opened squares) from  $750\text{ cm}^{-1}$  as a function of the Raman linewidth  $\Gamma$  and the correlation length  $L$  (closed triangles) as a function of Raman linewidth  $\Gamma$  are displayed in (a) and (b), respectively. The Raman linewidth performance follows a trend which displays a maximum for M12, a value at which a maximum disorder is supposed to be expected in a random system.

correlation length behavior for the AlGaIn-like  $A_1$  (LO) phonon. The Raman shift  $\Delta\omega$  (opened squares) from  $750\text{ cm}^{-1}$  as a function of the Raman linewidth  $\Gamma$  and the correlation length  $L$  (closed triangles) as a function of Raman linewidth  $\Gamma$  are displayed in Figs. 4(a) and 4(b), respectively. The Raman linewidth performance follows a trend which displays a maximum for M12, a value at which a maximum disorder is supposed to be expected in a random system. Thus, our data suggest that the distribution of the InAlGaIn/GaN alloy constituents for the higher Al to In ratio is random. It is also shown that the agreement between the experimental and calculated linewidth and shift of AlGaIn-like  $A_1$  (LO) phonon signal is reasonable. In the case of a single crystal with very low non-stoichiometry, the correlation length will tend towards infinity. In this study, the correlation length was discussed as a fitting parameter to investigate the localized region of AlGaIn-like  $A_1$  (LO) phonon. The value of  $L$  increases with decreasing Al to In ratio, which indicates that the phonon extended region becomes larger. Hence, we have demonstrated experimentally that the  $L$  value is a

very appropriate parameter accounting for the disorder of InAlGaIn/GaN heterostructures.

#### IV. CONCLUSIONS

In summary, the asymmetric behavior of the AlGaIn-like  $A_1$  (LO) Raman linewidth of InAlGaIn/GaN heterostructures was attributed to the activation of phonons of  $q \cong 0$  arising from the disordered state of the alloys. Thus, the microscopic nature of the InAlGaIn/GaN disorder is discussed by investigating the compositional dependence of AlGaIn-like  $A_1$  (LO) mode and is able to be quantitatively explained in terms of the SC model. X-ray diffraction supports the results; the correlation length  $L$  decreases as the increasing of Al to In ratio corresponding to the absence of the long-range order in the alloy.

#### ACKNOWLEDGMENTS

The authors would like to acknowledge the support of the National Science Council Project Nos. NSC98-2221-E-272-005 and NSC99-2112-M-236-001-MY3.

- <sup>1</sup>H. Hirayama, *J. Appl. Phys.* **97**, 091101 (2005).
- <sup>2</sup>S. Nagarajan, T. S. Oh, M. S. Kumar, C. H. Hong, and E. K. Suh, *Jpn. J. Appl. Phys., Part 1* **47**, 4413 (2008).
- <sup>3</sup>C. H. Chen, Y. F. Chen, Z. H. Lan, L. C. Chen, K. H. Chen, H. X. Jiang, and J. Y. Lin, *Appl. Phys. Lett.* **84**, 1480 (2004).
- <sup>4</sup>V. Yu. Davydov, I. N. Goncharuk, A. N. Smirnov, A. E. Nikolaev, W. V. Lundin, A. S. Usikov, A. A. Klochikhin, J. Aderhold, J. Graul, O. Semchinova, and H. Harima, *Phys. Rev. B* **65**, 125203 (2002).
- <sup>5</sup>A. Cros, A. Cantarero, N. T. Pelekanos, G. J. Pomeroy, and M. Kuball, *Phys. Status Solidi B* **243**, 1674 (2006).
- <sup>6</sup>T. Yu, Y. Pan, Z. Yang, K. Xu, and G. Zhang, *J. Cryst. Growth* **298**, 211 (2007).
- <sup>7</sup>J. Wu, J. Li, G. Cong, H. Wei, P. Zhang, W. Hu, X. Liu, Q. Zhu, Z. Wang, Q. Jia, and L. Guo, *Nanotechnology* **17**, 1251 (2006).
- <sup>8</sup>J. B. Wang, H. M. Zhong, Z. F. Li, and W. Lu, *J. Appl. Phys.* **97**, 086105 (2005).
- <sup>9</sup>P. Parayanthal and F. H. Pollak, *Phys. Rev. Lett.* **52**, 1822 (1984).
- <sup>10</sup>L. Y. Lin, C. W. Chang, W. H. Chen, and Y. F. Chen, *Phys. Rev. B* **69**, 075204 (2004).
- <sup>11</sup>I. H. Campbell and P. M. Fauchet, *Solid State Commun.* **58**, 739 (1986).
- <sup>12</sup>K. K. Tiong, P. M. Amirtharaj, F. H. Pollak, and D. E. Aspnes, *Appl. Phys. Lett.* **44**, 122 (1984).
- <sup>13</sup>S. Y. Hu, Y. C. Lee, Z. C. Feng, and S. H. Yang, *J. Lumin.* **132**, 1037 (2012).
- <sup>14</sup>H. X. Jiang and J. Y. Lin, *Opt. Electron. Rev.* **10**, 271 (2002).